# Geometric and electronic structures of new endohedral fullerenes: $\mathbf{E u} @ \mathbf{C}_{72}$ 

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#### Abstract

The geometric and electronic structures of rare earth metallofullerenes $\mathrm{Eu} @ \mathrm{C}_{72}$ were investigated using density functional theory (DFT) within a generalized gradient approximation (GGA). The geometric optimization revealed that the most favorable endohedral site for Eu is off-center along the $\mathrm{C}_{2}$ axis on the $\sigma_{\mathrm{v}}$ plane pointing to the $(5,5)$ bond at the fusion of two pentagons. Calculations for electronic structures show that two 6 s electrons in Eu transfer to the lowest-unoccupied-molecular orbitals of $\mathrm{C}_{72}$ while 4f electrons remain in Eu.


Keywords Metallofullerenes•Eu@C $\mathrm{C}_{72}$.
Density functional theory

## Introduction

Fullerenes are closed-shell molecular cages made up of carbon atoms that are connected together to form five- and

[^0]six-membered rings [1-3]. Most fullerenes obey the socalled isolated pentagon rule (IPR), which means that each pentagon is surrounded by five hexagons. In general, this arrangement is the most stable because neighboring pentagons increase the strain in the system and disrupt the ideal carbon-carbon bonding pattern. The IPR has proved particularly valuable in unraveling the cage structures of higher fullerenes and metallofullerenes [4-9]. Fullerenes, if they violate this rule, are believed to be unstable and difficult to isolate. However, deviations from the IPR are have sometimes been found in some experimental studies of endohedral fullerenes such as $\mathrm{Sc}_{2} @ \mathrm{C}_{66}$ [10], $\mathrm{Sc}_{3} \mathrm{C} @ \mathrm{C}_{68}$ [11, 12], $\mathrm{Tb}_{3} \mathrm{~N} @ \mathrm{C}_{84}$ [13], etc. Theoretical investigations of $\mathrm{C}_{72}$ [14], $\mathrm{Ca} @ \mathrm{C}_{72}$ [15] and $\mathrm{Mg} @ \mathrm{C}_{72}$ [16] also suggest that non-IPR cage structures are more stable than the IPRsatisfying structure. Recently, as the first proven structure of the $\mathrm{C}_{72}$ family, $\mathrm{La} @ \mathrm{C}_{72}$ was isolated and characterized by Wakahara et al. [17]. Experimental and DFT results suggested that $\mathrm{La} @ \mathrm{C}_{72}$ has a non-IPR $\mathrm{C}_{2}$ cage and that the encapsulated La atom locates near the adjacent pentagons $[17,18]$. These studies revealed that the inner metal atom(s) is(are) bonded to the fused pentagons, and electrons transfer from the inner metal core to the carbon cage, which change the electronic structures of the carbon cage, leading to the stabilization of the non-IPR fullerenes and conferring their special properties.

The rare earth element europium (Eu) has the electronic configuration $(\mathrm{Xe})^{54}(4 \mathrm{f})^{7}(6 \mathrm{~s})^{2}$ and its classical chemistry involves the +2 and +3 redox states. A series of compounds activated by the Eu ion have been studied for their practical application as phosphorescent materials [19, 20]. The interaction between the Eu ion and carbon cages might influence the spectrum and lead to a new matrix of luminescence. Since the successful extraction of $\mathrm{Eu} @ \mathrm{C}_{74}$ in 1998 [21], europium endohedral metallofullerenes have attracted special attention because of their potentially
luminescence properties, which cannot be expected from empty fullerenes [22, 23]. Very recently, endohedral fullerence Eu@C $\mathrm{C}_{72}$ was successfully synthesized and separated by Bucher et al. [24, 25]. Two isomers were separated by HPLC and identified by LDI-TOF mass spectrometry; however, detailed information on their structural and electronic properties is still unknown. In this study, density functional theory (DFT) calculations were performed on $\mathrm{Eu} @ \mathrm{C}_{72}$ in order to better understand the experiments of Bucher et al. and to explore the structural and electronic properties of the cluster.

## Methods

All DFT computations of $\mathrm{C}_{72}$ and $\mathrm{Eu} @ \mathrm{C}_{72}$ were performed using Dmol $^{3}$ code [26-28] with Perdew-Wang's 1991 generalized gradient approximation (GGA) function (PW91) [29]. An all-electron double-numerical basis set with polarization functions (DNP) was applied for all empty and endohedral cages. It is known that relativistic effects play an important role in the chemical and physical properties of molecules containing heavier elements such as europium. To take into account relativistic effects, we chose the all-electron scalar relativistic method for all calculations, utilizing the Douglas-Kroll-Hess (DKH) Hamiltonian [30], which is the most accurate approach available in the Dmol ${ }^{3}$ package. A self-consistent field procedure was carried out, with a convergence criterion of $10^{-6}$ a.u. on energy and electron density, and geometry optimizations were performed under the symmetry constraint with a convergence criterion of $10^{-3}$ a.u. on gradient, $10^{-3}$ a.u. on displacement, and $10^{-5}$ a.u. on energy.

According to the research of Nagase et al. [31], three cages were considered in this paper, namely, one IPR cage with $\mathrm{D}_{6 \mathrm{~d}}$ symmetry and two non-IPR isomers containing one fused pentagon with $\mathrm{C}_{2 \mathrm{v}}$ and $\mathrm{C}_{2}$ symmetry, respectively. By placing the Eu atom at different sites among these isomers, nine configurations of $\mathrm{Eu} @ \mathrm{C}_{72}$ were obtained. For the IPR cage, two sites were given: (1) in the center $\left(\mathrm{Eu} @ \mathrm{C}_{72}-1\right)$; (2) off-center on the $\sigma_{\mathrm{d}}$ plane ( $\mathrm{Eu} @ \mathrm{C}_{72}-2$ ). For the $\mathrm{C}_{2 \mathrm{v}}$ cage, four different positions of the Eu atom were considered: (3) in the center ( $\mathrm{Eu} @ \mathrm{C}_{72}-3$ ); (4) offcenter site along the $\mathrm{C}_{2}$ axis on the $\sigma_{\mathrm{v}}$ plane pointing to the $(5,5)$ bond, the fusion of two pentagons (Eu@C $\mathrm{C}_{2}-4$ ); (5) off-center site along the $\mathrm{C}_{2}$ axis on the $\sigma_{\mathrm{v}}$ plane pointing to the $(6,6)$ bond, the fusion of two hexagons $\left(\mathrm{Eu} @ \mathrm{C}_{72}-5\right)$; (6) off-center site along the axis vertical to the $\mathrm{C}_{2}$ axis under a pentagon center on the $\sigma_{v}$ plane $\left(\mathrm{Eu} @ \mathrm{C}_{72}-6\right)$. The positions of the Eu atom in the $\mathrm{C}_{2}$ cage are: (7) in the center $\left(\mathrm{Eu} @ \mathrm{C}_{72}-7\right)$; (8) off-center site along the $\mathrm{C}_{2}$ axis pointing to the $(5,5)$ bond, the fusion of two pentagons $\left(\mathrm{Eu} @ \mathrm{C}_{72}-8\right)$ and (9) off-center site along the $\mathrm{C}_{2}$ axis pointing to the (6, 6) bond, the fusion of two hexagons ( $\mathrm{Eu} @ \mathrm{C}_{72}-9$ ). Figure 1 shows the nine possible structures of $\mathrm{Eu} @ \mathrm{C}_{72}$.

## Results and discussion

## Relative stability of $\mathrm{C}_{72}$ isomers

Considering the neutral order first, $\mathrm{C}_{2 \mathrm{v}} \mathrm{C}_{72}$ with one pentagon-pentagon fusion has the lowest energy, followed by the $\mathrm{D}_{6 \mathrm{~d}}$ IPR isomer ( $65.39 \mathrm{~kJ} \mathrm{~mol}^{-1}$ higher in energy) and $\mathrm{C}_{2} \mathrm{C}_{72}$ ( $79.05 \mathrm{~kJ} \mathrm{~mol}^{-1}$ higher in energy) (Table 1). The isomeric stability is in quite good agreement with the


Fig. 1 Nine possible exohedral configurations for $\mathrm{Eu} @ \mathrm{C}_{72}$. Green ball Europium (Eu), gray balls the cage

Table 1 Number of pentagon-pentagon fusions $\left(\mathrm{N}_{\mathrm{pp}}\right)$, and the PW91/ DNP relative energies of $\mathrm{C}_{72}$ and $\mathrm{C}_{72}{ }^{2-}$

| Isomer | $\mathrm{N}_{\mathrm{pp}}$ | $\mathrm{C}_{72}$ <br> $\mathrm{E}_{\text {rel }}\left(\mathrm{kJ} \mathrm{mol}^{-1}\right)$ | $\mathrm{C}_{72}{ }^{2-}$ <br> $\mathrm{E}_{\text {rel }}\left(\mathrm{kJ} \mathrm{mol}^{-1}\right)$ |
| :--- | :--- | :--- | :--- |
| $\mathrm{D}_{6 \mathrm{~d}}$ | 0 | 65.39 | 109.25 |
| $\mathrm{C}_{2 \mathrm{v}}$ | 1 | 0 | 0 |
| $\mathrm{C}_{2}$ | 1 | 79.05 | 13.13 |

result in [14], in which Becke's three parameter function combined with the non-local Lee-Yang-Parr correlation function (B3LYP) in the standard 3-21G basis set was used. The screening of a metallofullerene structure usually begins by considering the charged empty cages; the magnitude of the negative charge is based on the expected electron donation from the encapsulated metal. Since the electronic structure of $\mathrm{Eu} @ \mathrm{C}_{72}$ can be described as $\mathrm{Eu}^{2+}$ $@ \mathrm{C}_{72}{ }^{2-}[24,25]$, the $\mathrm{C}_{72}{ }^{2-}$ isomers were also computed. While the stability order of $\mathrm{C}_{72}$ dianions changes, the stablest isomer is still the non-IPR isomer $\mathrm{C}_{2 \mathrm{v}}$, followed by the $\mathrm{C}_{2}$ (relative energy $13.13 \mathrm{~kJ} \mathrm{~mol}^{-1}$ ). The IPR isomer becomes the most unstable, with a markedly higher relative energy (109. $\mathrm{kJ} \mathrm{mol}^{-1}$ ).

## Relative stability of $\mathrm{Eu} @ \mathrm{C}_{72}$ isomers

Table 2 presents the HOMO, LUMO, $\mathrm{E}_{\text {gap }}$, and binding energy (BE) of all nine isomers shown in Fig. 1. Commonly, the thermodynamic stability of the fullerene cage is determined by the BE , which is defined as the absolute value of the difference between the total energy of the molecule and the energy sum of all free atoms constituting the molecule. Obviously, Eu@ $\mathrm{C}_{72}-4$ has the lowest BE of -547.49 eV , followed by $\mathrm{Eu} @ \mathrm{C}_{72}-7$. So, the most stable site for the Eu atom is the off-center site along the $\mathrm{C}_{2}$ axis on the $\sigma_{\mathrm{v}}$ plane pointing to the $(5,5)$ bond, the fusion of

Table 2 The highest-occupied-molecular orbital (HOMO), lowest-unoccupied-molecular orbital (LUMO), $\mathrm{E}_{\text {gap }}$, and binding energy (BE) of the nine $\mathrm{Eu} @ \mathrm{C}_{72}$ isomers (in eV )

| Isomers | BE | HOMO | LUMO | Egap |
| :--- | :--- | :--- | :--- | :--- |
| $\mathrm{Eu@C}_{72}-1$ | -544.10 | -4.88 | -4.72 | 0.16 |
| $\mathrm{Eu} @ \mathrm{C}_{72}-2$ | -545.93 | -4.98 | -4.73 | 0.25 |
| $\mathrm{Eu} @ \mathrm{C}_{72}-3$ | -545.23 | -5.14 | -5.02 | 0.12 |
| $\mathrm{Eu} @ \mathrm{C}_{72}-4$ | -547.49 | -5.18 | -4.73 | 0.45 |
| $\mathrm{Eu} @ \mathrm{C}_{72}-5$ | -546.28 | -5.31 | -4.87 | 0.44 |
| $\mathrm{Eu} @ \mathrm{C}_{72}-6$ | -546.71 | -5.14 | -4.94 | 0.20 |
| $\mathrm{Eu} @ \mathrm{C}_{72}-7$ | -547.31 | -5.45 | -4.81 | 0.64 |
| $\mathrm{Eu} @ \mathrm{C}_{72}-8$ | -547.31 | -5.45 | -4.81 | 0.64 |
| $\mathrm{Eu} @ \mathrm{C}_{72}-9$ | -546.29 | -5.32 | -5.00 | 0.32 |



Fig. 2 Electron spin density of $\mathrm{Eu} @ \mathrm{C}_{72}-4$, the illustrated hypersurface corresponding to a value of $0.01 \mathrm{e} \mathrm{au}^{-3}$
two pentagons in the $\mathrm{C}_{2 \mathrm{v}}$ cage. One more thing worth noting is that the Eu atom in the center of the $\mathrm{C}_{2}$ cage moves to the Eu site in $\mathrm{Eu} @ \mathrm{C}_{72}-7$ during optimization. $\mathrm{Eu} @ \mathrm{C}_{72}-7$ and $\mathrm{Eu} @ \mathrm{C}_{72}-8$ converge to the same structure.

We also compared the $\mathrm{E}_{\mathrm{gap}}$ values of all the considered structures, because the $\mathrm{E}_{\text {gap }}$ value is often correlated with kinetic stability. The $\mathrm{E}_{\text {gap }}$ values calculated for the hollow cages are $1.45,0.66$ and 0.30 eV for $\mathrm{D}_{6 \mathrm{~d}}, \mathrm{C}_{2 \mathrm{v}}$ and $\mathrm{C}_{2} \mathrm{C}_{72}$, respectively. It can be seen from Table 2 that, among the nine structures, $\mathrm{Eu} @ \mathrm{C}_{72}-7$ has the largest $\mathrm{E}_{\text {gap }}$, so it is expected to be the most stable kinetically. On the other hand, $\mathrm{Eu} @ \mathrm{C}_{72}-3$ has the smallest $\mathrm{E}_{\text {gap }}$ of 0.12 eV and is therefore regarded as the most unstable kinetically. Moreover, except for the $\mathrm{C}_{2}$ isomer, the $\mathrm{E}_{\text {gap }}$ values of the $\mathrm{D}_{6 \mathrm{~d}}$ and $\mathrm{C}_{2 \mathrm{v}}$ cages are slightly reduced by encaging the Eu atom.


Fig. 3 Electronic density in the plane passing through the Eu ion in $\mathrm{Eu} @ \mathrm{C}_{72}-4$


Fig. 4 Electrostatic potential map for $\mathrm{C}_{72}{ }^{2-}$ in the plane passing through the $\sigma_{\mathrm{v}}$ plane

Here, we focus on the properties of $\mathrm{Eu} @ \mathrm{C}_{72}-4$, which has the lowest BE. For $\mathrm{Eu} @ \mathrm{C}_{72}-4$, the Eu atom is situated away from the center of $\mathrm{C}_{72}$ by about $1.40 \AA$. The $\mathrm{C}-\mathrm{C}$ bond lengths range from $1.39 \AA$ to $1.49 \AA$, which differ slightly from those of $\mathrm{C}_{72}$, and the shortest $\mathrm{Eu}-\mathrm{C}$ bond length is $2.57 \AA$. In $\mathrm{Eu} @ \mathrm{C}_{60}$ and $\mathrm{Eu} @ \mathrm{C}_{82}$, Eu ions are positioned slightly away from the center, $1.20 \AA$ [32] for the former and $1.40 \AA$ [33] for the latter. The Eu-C bond lengths in $\mathrm{Eu} @ \mathrm{C}_{72}$ are $0.22 \AA$ shorter than that in $\left(\mathrm{C}_{5} \mathrm{Me}_{5}\right)_{2} \mathrm{Eu}$ [34].

The total spin multiplicity of the ground state is $\mathrm{S}=8$. The spin densities on Eu and $\mathrm{C}_{72}$ are 6.964 and 0.036 , respectively. The spin multiplicities of $S=2,4$ and 6 ,are $386.44,187.80$ and $41.85 \mathrm{~kJ} \mathrm{~mol}^{-1}$ higher respectively, than that of $S=8$. Figure 2 shows the spatial distribution of spin density for $\mathrm{Eu} @ \mathrm{C}_{72}-4$. It can be seen from Fig. 2 that the spin densities are distributed over the Eu atom.

Mulliken population analysis shows that the resulting electronic configuration of Eu is $6 \mathrm{~s}^{0.14} 4 \mathrm{f}^{6.972} 5 \mathrm{~d}^{0.771}$ and the Mulliken charge on Eu is +0.973 e . The occupation of 0.771 e on the Eu 5d orbitals is ascribed to a contribution of


Fig. 6 The total density of state (TDOS) of $\mathrm{C}_{72}, \mathrm{Eu} @ \mathrm{C}_{72}-4$, and partial density of state (PDOS) of Eu in $\mathrm{Eu} @ \mathrm{C}_{72}-4$. Dotted line Fermi energy in units of eV
back-donation of the occupied orbitals of the fullerene cage [35-37]. After this correction, the charge transfer from Eu to $\mathrm{C}_{72}$ is about 1.744 e , which is close to the experimental result. Therefore, seven electrons spin on the Eu ion and two valance electrons transfer onto the $\mathrm{C}_{72}$ cage. The electronic structure of $\mathrm{Eu} @ \mathrm{C}_{72}$ can be described as $\mathrm{Eu}^{2+} \mathrm{C}_{72}{ }^{2-}$.

In addition to the charge transfer, orbital hybridization between the guest transition-metal/lanthanide atoms and fullerene is another common characteristic of monometallofullerenes [38-41], and we find this to be so in $\mathrm{Eu} @ \mathrm{C}_{72}$. Figure 3 displays the electron density in a plane containing the Eu atom. Minor electron accumulation leads to weak covalent bond formation between Eu and proximal C atoms. Therefore, the interaction between Eu and $\mathrm{C}_{72}$ is primarily ionic with a weak covalent character.

Electrostatic interactions play an important role in determining the position of the positively charged Eu ion. To provide insight into these interactions, the electrostatic potential map inside the $\mathrm{C}_{2 \mathrm{v}} \mathrm{C}_{72}{ }^{2-}$ anion was calculated. Figure 4 shows the electrostatic potential map for $\mathrm{C}_{72}{ }^{2-}$ in the plane passing through $\sigma_{\mathrm{v}}$. It can be seen that the


Fig. 5 Isosurfaces (the isovalue is 0.03 a.u.) of: a HOMO-1, b HOMO, c LUMO, d LUMO +1 in Eu@C $\mathrm{C}_{72}$. Yellow and blue are used to indicate the positive and negative sign of the wavefunction,

c)
respectively (for interpretation of the references to the color in this figure legend, the reader is referred to the web version of this article)
electrostatic potential become more negative upon going from the area near the $\mathrm{C}-\mathrm{C}$ double bond to the area near the pentagonal ring. This confirms our previous result that Eu@ $\mathrm{C}_{72}-4$, in which Eu ion occupies a more negative site, is the most stable structure.

The orbital levels of $\mathrm{C}_{72}$ and $\mathrm{Eu} @ \mathrm{C}_{72}-4$ are shown in Fig. 5. We label the highest carbon-derived occupied state (C-HOMO) and the lowest carbon-derived unoccupied state (C-LUMO) as those occupied (or empty) states nearest to the Fermi level that show barely any weight of the Eu atom [42]. The HOMO-1 and LUMO of Eu@C 72 should be the C-HOMO and C-LUMO, respectively. The energy difference between the C-HOMO and C-LUMO states is 0.56 eV for $\mathrm{Eu} @ \mathrm{C}_{72}, 0.1 \mathrm{eV}$ smaller than that for $\mathrm{C}_{2 \mathrm{v}} \mathrm{C}_{72}$.

To better understand the electronic structure of this endohedral fullerene, the total density of state (TDOS) of $\mathrm{Eu} @ \mathrm{C}_{72}-4$ is plotted in Fig. 6, together with the TDOS for $\mathrm{C}_{72}$ and the partial density of state (PDOS) of Eu in $\mathrm{Eu} @ \mathrm{C}_{72}-4$. The DOS was obtained by a Lorentzian extension of the discrete energy levels, with weights being the orbital populations in the levels, and a summation over them. The broadening width parameter chosen was 0.20 eV and the Fermi level $\left(\mathrm{E}_{\mathrm{f}}\right)$ was taken as zero. We explored the hybridized orbitals by analyzing TDOS of $\mathrm{Eu} @ \mathrm{C}_{72}-4$ and PDOS of Eu in $\mathrm{Eu} @ \mathrm{C}_{72}-4$. The figure shows that there is no peak of PDOS for Eu in $\mathrm{Eu} @ \mathrm{C}_{72}-4$ near the $\mathrm{E}_{\mathrm{f}}$, while there are some small peaks appearing in the energy range from -10 to -1 eV , indicating some weak hybridization between the atomic orbitals of Eu and those of the carbons of cage $\mathrm{C}_{72}$.

Finally, to gain insight into the reactivity properties of $\mathrm{Eu} @ \mathrm{C}_{72}-4$ compared with pristine $\mathrm{C}_{2 \mathrm{v}} \mathrm{C}_{72}$, the vertical ionization potentials (VIPs) and the vertical electron affinities (VEAs) of $\mathrm{C}_{2 \mathrm{v}} \mathrm{C}_{72}$ and $\mathrm{Eu} @ \mathrm{C}_{72}-4$ were calculated. The VIP is the energy difference between positively charged and neutral clusters. The VEA is evaluated by adding one electron to the neutral cluster in its equilibrium geometry and taking the difference between their total energies. The calculated VIP and VEA are 6.56 and 2.88 eV for $\mathrm{C}_{2 \mathrm{v}} \mathrm{C}_{72}$, while the calculated values of VIP and VEA for $\mathrm{Eu} @ \mathrm{C}_{72}$ are 6.62 and 3.28 eV , respectively. Both the VIP and VEA are slightly higher than that of pristine $\mathrm{C}_{72}$, indicating that $\mathrm{Eu} @ \mathrm{C}_{72}$ has a slightly stronger ability to capture electrons and has more difficultly in losing electrons compared with pristine $\mathrm{C}_{72}$. We hope further experiments will confirm our calculations.

## Conclusions

All-electron relativistic DFT calculations were performed on nine possible optimized geometries of $\mathrm{Eu} @ \mathrm{C}_{72}$. Among the possible optimized geometries of $\mathrm{Eu} @ \mathrm{C}_{72}$, the most
favorable endohedral site of Eu is an off-center site along the $\mathrm{C}_{2}$ axis on the $\sigma_{\mathrm{v}}$ plane pointing to the $(5,5)$ bond in the $\mathrm{C}_{2 \mathrm{v}}$ cage. Calculations for electronic structures show that the spin densities are distributed over the Eu ion. Two valance electrons of Eu transfer onto the $\mathrm{C}_{72}$ cage, which is in quite good agreement with experimental results. In addition, the Eu ion has little distribution to DOS around $\mathrm{E}_{\mathrm{f}}$.

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